

§ Eigenvalues

§ 01 The wave equation on a string

§ 02 The heat equation

§ 03 The Dirichlet eigenvalue problem

§ 04 The Neumann eigenvalue problem

§ 05 Eigenvalues of Laplace-Beltrami operator on S^2

Using spherical coordinates (θ, ϕ) , the Laplace-Beltrami operator on S^2 is :

$$\Delta_{S^2} = \frac{1}{\sin \theta} \frac{\partial}{\partial \theta} (\sin \theta \frac{\partial}{\partial \theta}) + \frac{1}{\sin^2 \theta} \frac{\partial^2}{\partial \phi^2} \dots (*)$$

We seek eigenfunctions $Y(\theta, \phi)$, and eigenvalues λ such that $\Delta_{S^2} Y = -\lambda Y$

Separation of variables

設 $Y(\theta, \phi) = \Theta(\theta)\Phi(\phi)$ 代入 (*) 得

$$\frac{1}{\sin \theta} \frac{\partial}{\partial \theta} (\sin \theta \frac{\partial \Theta}{\partial \theta}) \Phi + \frac{1}{\sin^2 \theta} \Theta \frac{\partial^2 \Phi}{\partial \phi^2} = -\lambda \Theta \Phi, \text{ 同除以 } \Theta \Phi, \text{ 再同乘以 } \sin^2 \theta$$

$$\frac{\sin \theta}{\Theta} \frac{d}{d\theta} (\sin \theta \frac{d\Theta}{d\theta}) + \lambda \sin^2 \theta + \frac{1}{\Phi} \frac{d^2 \Phi}{d\phi^2} = 0$$

設 $\frac{1}{\Phi} \frac{d^2 \Phi}{d\phi^2} = -m^2$, 則

$$\left\{ \begin{array}{l} \frac{d^2 \Phi}{d\phi^2} + m^2 \Phi = 0 \dots (1) \\ \frac{\sin \theta}{\Theta} \frac{d}{d\theta} (\sin \theta \frac{d\Theta}{d\theta}) + (\lambda \sin^2 \theta - m^2) \Theta = 0 \dots (2) \end{array} \right.$$

(1) 式是方位角方程式, (2)式是極角方程式。

由(1) $\Phi(\phi) = e^{im\phi}$, 其中 m 是整數(因為 BC $\Phi(\phi + 2\pi) = \Phi$)

(2) 的部分, 令 $x = \cos \theta$, $\sin \theta = \sqrt{1-x^2}$, $\frac{d}{d\theta} = -\sin \theta \frac{d}{dx} = -\sqrt{1-x^2} \frac{d}{dx} \dots (3)$

設 $P(x) = \Theta(\theta)$ 代入(3) $\frac{d\Theta}{d\theta} = -\sqrt{1-x^2} \frac{dP}{dx}$, $\sin \theta \frac{d\Theta}{d\theta} = -(1-x^2) \frac{dP}{dx}$

計算 $\frac{d}{d\theta} (\sin \theta \frac{d\Theta}{d\theta}) = \frac{d}{d\theta} (- (1-x^2) \frac{dP}{dx}) = \sqrt{1-x^2} \frac{d}{dx} ((1-x^2) \frac{dP}{dx})$

$\sin \theta \frac{d}{d\theta} (\sin \theta \frac{d\Theta}{d\theta}) = (1-x^2) \frac{d}{dx} ((1-x^2) \frac{dP}{dx})$

(2)式變成 $(1-x^2) \frac{d}{dx} ((1-x^2) \frac{dP}{dx}) + (\lambda(1-x^2) - m^2) P = 0$ 同處以 $(1-x^2)$ 後展開得

$$(1-x^2) \frac{d^2 P}{dx^2} - 2x \frac{dP}{dx} + (\lambda - \frac{m^2}{1-x^2}) P = 0$$

得到 associated Legendre equation :

$$\frac{d}{dx} \left((1-x^2) \frac{dP}{dx} \right) + \left(\lambda - \frac{m^2}{1-x^2} \right) P = 0$$

Eigenvalue $\lambda = l(l+1)$ 由 associated Legendre equation 解的存在性決定。

最後得到 eigenfunctions $Y_l^m(\theta, \phi) \propto P_l^m(\cos \theta) e^{im\phi}$, $\lambda_l = l(l+1)$

後記

1. 完整的解為 $Y_l^m(\theta, \phi) = \sqrt{\frac{2l + (l-m)!}{4\pi(1+m)!}} P_l^m(\cos \theta) e^{im\phi}$
2. 經典物理中，重力與靜電場滿足 $\Delta \Phi = 0$ ，當問題約束在球面上時就變成 $\Delta_{S^2} \Phi = 0$ ；例如球面上的熱擴散方程 $\frac{\partial u}{\partial t} = \Delta_{S^2} u$
3. 陀螺運動方程的(1)對稱性 (2)角動量量子化 (3)球面幾何與 Laplace-Beltrami operator 有深層的關係。
4. 陀螺的轉動 vs SO(3)，Laplace-Beltrami operator 是 SO(3)的 Casimir 算子！？