

Consider the heat equation :

$$\begin{cases} u_t = u_{xx}, (x,t) \in (-\infty, \infty) \times (0, \infty) \\ u(x,0) = f(x), x \in R \end{cases}, \text{ where } f(x) \text{ is continuous and } 0 \leq f(x) \leq 1$$

(a) Show that the equation has a solution u which satisfies $\lim_{t \rightarrow 0^+} u(x,t) = f(x)$ for $x \in R$

(b) Show that there is an $f(x)$ with $0 \leq f(x) \leq 1$ such that the equation has a solution u satisfying $\limsup_{t \rightarrow \infty} u(0,t) = 1$ and $\liminf_{t \rightarrow \infty} u(0,t) = 0$

The fundamental solution is $G(x,t) = \frac{1}{\sqrt{4\pi t}} \exp(-\frac{x^2}{4t})$ and

$$u(x,t) = \int_{-\infty}^{\infty} f(y) \cdot \frac{1}{\sqrt{4\pi t}} \exp(-\frac{(x-y)^2}{4t}) dy \text{ is a particular solution.}$$

(a)是要求證明 solution $u(x,t)$ 在 $t \rightarrow 0^+$ 時收斂到 initial condition $f(x)$
對於無限區間的熱方程，我們使用傅立葉變換來求解。

定義傅立葉變換： $\hat{u}(k,t) = \int_{-\infty}^{\infty} u(x,t) e^{-ikx} dx$

對熱方程兩邊取傅立葉變換，利用導數的性質： $\frac{\partial}{\partial t} \hat{u}(k,t) = -k^2 \hat{u}(k,t)$

這是一個一階常微分方程，通解為： $\hat{u}(k,t) = \hat{f}(k) e^{-k^2 t}$

其中 $\hat{f}(k) = \int_{-\infty}^{\infty} f(x) e^{-ikx} dx$ 是 $f(x)$ 的 Fourier transform

由反傅立葉變換 $u(x,t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} \hat{f}(k) e^{-k^2 t} e^{ikx} dx$

這是 initial function $f(x)$ 與 Gaussian kernel 的卷積： $u(x,t) = \int_{-\infty}^{\infty} f(y) \frac{1}{\sqrt{4\pi t}} e^{-\frac{(x-y)^2}{4t}} dy$

我們希望證明 $\lim_{t \rightarrow 0^+} u(x,t) = f(x)$ ，這相當於證明高斯核 $\frac{1}{\sqrt{4\pi t}} e^{-\frac{(x-y)^2}{4t}}$ 在 $t \rightarrow 0^+$ 時收

斂到 Dirac δ -function $\delta(x-y)$ 。

如果 $f(x)$ 是連續函數，則卷積(convolution)積分 $\lim_{t \rightarrow 0^+} \int_{-\infty}^{\infty} f(y) \frac{1}{\sqrt{4\pi t}} e^{-\frac{(x-y)^2}{4t}} dy = f(x)$

因此，當 $t \rightarrow 0^+$ 時 $u(x,t)$ 收斂到 initial condition $f(x)$ ，即 $\lim_{t \rightarrow 0^+} u(x,t) = f(x)$ 。

The fundamental solution to the heat equation $u_t = u_{xx}$ is given by :

$$K(x,t) = \frac{1}{\sqrt{4\pi t}} e^{-\frac{x^2}{4t}} \circ$$

This function satisfies the heat equation and has the property that : $\lim_{t \rightarrow 0^+} K(x,t) = \delta(x)$,

where $\delta(x)$ is the Dirac delta function \circ

The solution $u(x,t)$ to the heat equation with the initial condition $u(x,0)=f(x)$ can be expressed as the convolution of $f(x)$ with the heat kernel $K(x,t)$:

$$u(x,t) = \int_{-\infty}^{\infty} K(x-y,t)f(y)dy \circ$$

Substituting the expression for $K(x,t)$, we get : $u(x,t) = \frac{1}{\sqrt{4\pi t}} \int_{-\infty}^{\infty} e^{-\frac{(x-y)^2}{4t}} f(y)dy$

As $t \rightarrow 0^+$, the heat kernel $K(x-y, t)$ becomes highly concentrated around $y=x$ \circ This means that the integral will be dominated by the values of $f(y)$ near $y=x$ \circ

Formally , we can use the fact that the heat kernel approximates the Dirac delta function

as $t \rightarrow 0^+$: $\lim_{t \rightarrow 0^+} K(x-y,t) = \delta(x-y)$

Therefore, the convolution integral becomes : $\lim_{t \rightarrow 0^+} u(x,t) = \int_{-\infty}^{\infty} \delta(x-y)f(y)dy = f(x)$

§ Let $K(x,t) = \frac{1}{\sqrt{4\pi t}} \exp(-\frac{x^2}{4t})$, $x \in R$, $t > 0$

$f(x) \in C(R)$, $f(x+1)=f(x)$ and $\int_0^1 f(x)dx = 3$ \circ Define $u(x,t) = \int_R K(x-y,t)f(y)dy$

(a) Show that $\lim_{(x,t) \rightarrow (x_0, 0^+)} u(x,t) = f(x_0)$ for each $x_0 \in R$

(b) Show that $\lim_{t \rightarrow \infty} u(0,t) = 3$

(a) $u(x,t) = \int_R K(x-y,t)f(y)dy$, where $K(x,t) = \frac{1}{\sqrt{4\pi t}} \exp(-\frac{x^2}{4t})$ is the heat

kernel \circ

The heat kernel $K(x,t)$ has the following properties :

1. Normalization : $\int_{\mathbb{R}} K(x,t)dx = 1$ for all $t > 0$
2. Delta function behavior : As $t \rightarrow 0^+$, $K(x,t)$ tends to the Dirac delta function $\delta(x)$. This means that for any continuous function $f(x)$,

$$\lim_{t \rightarrow 0^+} \int_{\mathbb{R}} K(x-y,t)f(y)dy = f(x)$$

The function $f(x)$ is periodic 1 , implies that $f(x)$ can be represented as a Fourier series : $f(x) = \sum_{n=-\infty}^{\infty} c_n e^{2\pi i n x}$, where c_n are the Fourier coefficients .

The Fourier transform of the heat kernel $K(x,t)$ is $\hat{K}(\xi,t) = \exp(-4\pi^2 \xi^2 t)$

This shows that the heat kernel acts as a low-pass filter in the frequency domain, attenuating high-frequency components as $t \rightarrow 0^+$

Using the properties of the heat kernel and the periodicity of $f(x)$, we can write:

$$u(x,t) = \int_{\mathbb{R}} K(x-y,t)f(y)dy.$$

As $t \rightarrow 0^+$, the heat kernel $K(x-y,t)$ becomes increasingly concentrated around $y = x$, and the integral effectively samples $f(y)$ at $y = x$. Therefore, we have:

$$\lim_{t \rightarrow 0^+} u(x,t) = f(x).$$

Since this holds for any $x \in \mathbb{R}$, it follows that:

$$\lim_{(x,t) \rightarrow (x_0,0^+)} u(x,t) = f(x_0).$$